

Wrenches

a generalized force acting on a rigid body consists of

a linear component (pure force), and
an angular component (pure moment)

acting at a point.

↑ this is important.

This generalized force is called a wrench and is a vector in \mathbb{R}^6 (for $SE(3)$):

$$F = \begin{Bmatrix} f \\ \tau \end{Bmatrix} \quad \begin{array}{l} f \in \mathbb{R}^3 \text{ linear part} \\ \tau \in \mathbb{R}^3 \text{ angular part} \end{array}$$

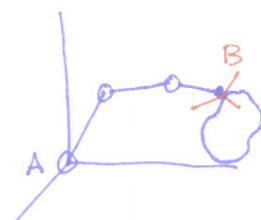
- a force/moment pair.
- values will depend on the coordinate frame.

Properties of wrenches:

1) combine w/ twists to define instantaneous work,

$$\delta W = \bar{F}_B \cdot \bar{\gamma}^B = (f \cdot v + \tau \cdot \omega)$$

↑ instantaneous body velocity (frame B)
applied wrench in frame B.
at point B.



• both operate at the same point.

2) over time (1) generates work: $W = \int_{t_1}^{t_2} \bar{F}_B \cdot \bar{\gamma}^B dt$

3) two wrenches are equivalent if they do the same work independent of $\bar{\gamma}, \bar{\gamma}$. (wrenches need not be at same point).

can use to go from body to spatial

$\cdot s = g^T$

1.5

$$F_A = \text{Ad}_{g_{BA}}^T F_B \quad \boxed{\text{_____}}$$



$$F_{\text{spatial}} = \text{Ad}_{g^{-1}}^T F_{\text{body}}$$

$$= (\text{Ad}_g^T)^{-1} F_{\text{body}}$$

$$F_{\text{body}} = \text{Ad}_g^T F_{\text{spatial}}$$

\uparrow this is flipped compared
to what we are used
to.

$$R^T - R^T P \underset{g}{\underset{\circ}{\gamma}} =$$

R

$$-R(P^T P)^\gamma$$

$$\xi^{\text{spatial}} - \text{Ad}_g \xi^{\text{body}}$$

$$F^{\text{sp}} \cdot \xi^{\text{sp}} = F^{\text{b}} \xi^{\text{b}}$$

$$F^{\text{sp}} \cdot \text{Ad}_g \xi^{\text{b}} - P^T \xi^{\text{b}}$$

$$\text{Ad}_g^T F^{\text{sp}} = F^{\text{b}} \xi^{\text{b}}$$

- recall, that this is as though new frame were rigidly attached to old frame.
This is NOT the same as just changing frame of reference of component description.

- wrenches can be added, but only if applied at the same point

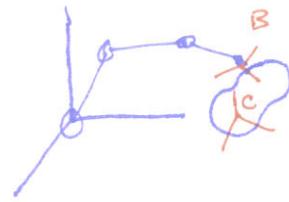
Theorem (Poinsot). Every collection of wrenches applied to a rigid body is equivalent to a force applied ^{along} a fixed axis plus a torque about the same axis

- dual theorem to Chasles' theorem.

$\cancel{F_B^A}$ - wrench in ~~B's frame~~ (at point B)

$\cancel{F_C^A}$ - wrench in ~~C's frame~~ (at point C)

both in A frame



ξ_B^A - twist at point B

ξ_C^A - twist at point C

equivalency means:

$$\begin{aligned} F_C \cdot \xi_C &= F_B \cdot \xi_B \\ &= F_B \cdot (\text{Ad}_{g_{BC}} \xi_C) \\ &= (\text{Ad}_{g_{BC}}^T F_B) \cdot \xi_C \end{aligned}$$

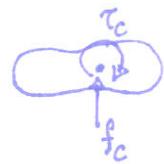
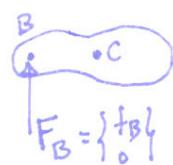
\Rightarrow

$$F_C = \text{Ad}_{g_{BC}}^T F_B$$

↑ transformation of wrench under coordinate frame transformation.

$$\begin{Bmatrix} f_C \\ \tau_C \end{Bmatrix} = \begin{bmatrix} R_{BC}^T & 0 \\ -R_{BC}^T \hat{p}_{BC} & R_{BC}^T \end{bmatrix} \begin{Bmatrix} f_B \\ \tau_B \end{Bmatrix} \quad \text{for } g_{BC} = (p_{BC}, R_{BC}).$$

you've seen this before:



pure free at B \rightarrow induces torque at C

Manipulator Dynamics

now, we know how to express the kinetic energy of a rigid body.

we already know how to express the potential energy of a rigid body.

how can this translate to a manipulator?

- it is a collection of rigid bodies w/ constraints
- can be described by the joint space
- if possible to define a lagrangian, equations of motion follow ~~natural~~ naturally.

Kinetic Energy:

- just need to sum up kinetic energy of each link.

let the center of mass of each link be $\mathbf{g}_{ci}(\theta)$

\Rightarrow

body velocity of link center-of-mass is

$$\dot{\xi}_{ci}^{\text{body}} = \mathbf{J}_{ci}^{\text{body}}(\theta)\dot{\theta}$$

\Rightarrow

kinetic energy of link is

$$T_i = \frac{1}{2} (\dot{\xi}_{ci}^{\text{body}})^T M_i (\dot{\xi}_{ci}^{\text{body}}) \quad \begin{matrix} \uparrow \\ \text{gen. inertia matrix of } i\text{-th link.} \end{matrix}$$

$$T_i(\theta, \dot{\theta}) = \frac{1}{2} \dot{\theta}^T (\mathbf{J}_{ci}^{\text{body}}(\theta))^T M_i \mathbf{J}_{ci}^{\text{body}}(\theta) \dot{\theta}$$

⇒

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{\theta}} - \frac{\partial L}{\partial \theta} = \sum_{j=1}^n M_{ij}(\theta) \ddot{\theta}_j + \sum_{j,k=1}^n \left(\frac{\partial M_{ij}}{\partial \theta_k} \dot{\theta}_j \dot{\theta}_k - \frac{1}{2} \frac{\partial M_{kj}}{\partial \theta_i} \dot{\theta}_k \dot{\theta}_j \right) + \frac{\partial V}{\partial \theta_i}(\theta) = \Upsilon_i$$

Typically written as:

$$\sum_{j=1}^n M_{ij}(\theta) \ddot{\theta}_j + \sum_{j,k=1}^n \Gamma_{ijk}(\theta) \dot{\theta}_j \dot{\theta}_k + \frac{\partial V}{\partial \theta_i}(\theta) = \Upsilon_i$$

where,

$$\Gamma_{ijk}(\theta) = \frac{1}{2} \left(\frac{\partial M_{ij}}{\partial \theta_k} + \frac{\partial M_{ik}}{\partial \theta_j} - \frac{\partial M_{kj}}{\partial \theta_i} \right)$$

\uparrow called Christoffel symbols
 corresponding to mass matrix
 $M(\theta)$.

\downarrow due to symmetry of M , $M = M^T$.

$$M(\theta) \ddot{\theta} + C(\theta, \dot{\theta}) \dot{\theta} + N(\theta, \dot{\theta}) = \tau$$

\uparrow inertial forces \uparrow Coriolis &
 centrifugal
 forces \uparrow potential &
 nonconservative
 forces \uparrow activation torques
 \uparrow Υ_i broken up into
 two parts.

• Second-order, vector differential equation.

$$(C(\theta, \dot{\theta})) \dot{\theta} = \sum_{k=1}^n \Gamma_{ijk}(\theta) \dot{\theta}_k = \frac{1}{2} \sum_{k=1}^n \left(\frac{\partial M_{ij}}{\partial \theta_k} + \frac{\partial M_{ik}}{\partial \theta_j} - \frac{\partial M_{kj}}{\partial \theta_i} \right) \dot{\theta}_k$$

$$N_i(\theta, \dot{\theta}) = \frac{\partial V}{\partial \theta_i} - \text{external forces}$$

\uparrow obtained from Υ_i

e.g. $\beta_i \dot{\theta}_i$

\uparrow friction at joint.

Total kinetic energy found by summing all of the link KE's:

$$T(\theta, \dot{\theta}) = \sum_{i=1}^n T_i(\theta, \dot{\theta}) = \frac{1}{2} \dot{\theta}^T M \dot{\theta}$$

$\uparrow M(\theta) \in \mathbb{R}^{n \times n}$ is called manipulator inertia matrix.

$$M(\theta) = \sum_{i=1}^n (J_{ci}^{\text{body}}(\theta))^T M_i J_{ci}^{\text{body}}(\theta)$$

Potential Energy:

- this one is simply a matter of substituting in the forward kinematics

$$\begin{aligned} V(g) &\rightarrow V_i(g_{ci}(\theta)) \\ \Rightarrow V(\theta) &= \sum_i^n V_i(g_{ci}(\theta)) = \sum_i^n V_i g_{ci}(\theta) \end{aligned}$$

Total Energy:

Complete Manipulator Lagrangian:

$$L(\theta, \dot{\theta}) = \sum_{i=1}^n (T_i(\theta, \dot{\theta}) - V_i(\theta)) = \frac{1}{2} \dot{\theta}^T M(\theta) \dot{\theta} - V(\theta)$$

Equations of Motion:

Once we have the manipulator Lagrangian, the rest is a piece of cake (kinda).

Just work out the ~~Lagrange~~ Lagrange's equations; but first ~~for~~ some preliminaries:

$$L(\theta, \dot{\theta}) = \frac{1}{2} \dot{\theta}^T M(\theta) \dot{\theta} - V(\theta)$$

$$= \frac{1}{2} \sum_{i,j=1}^n M_{ij}(\theta) \dot{\theta}_i \dot{\theta}_j - V(\theta)$$

Want to find

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{\theta}_i} - \frac{\partial L}{\partial \theta_i} = \tau_i$$

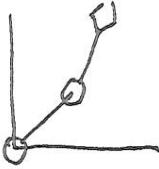
↑ activation torque & other nonconservative ~~forces~~, generalized forces acting on i^{th} joint.

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{\theta}_i} = \frac{d}{dt} \left(\sum_{j=1}^n M_{ij} \dot{\theta}_j \right) = \sum_{j=1}^n (M_{ij} \ddot{\theta}_j + \dot{M}_{ij} \dot{\theta}_j)$$

$$= \sum_{j=1}^n (M_{ij} \ddot{\theta}_j + \sum_{k=1}^n \frac{\partial M_{ij}}{\partial \theta_k} \dot{\theta}_k \dot{\theta}_j)$$

$$\frac{\partial L}{\partial \theta_i} = \frac{1}{2} \sum_{j,k=1}^n \frac{\partial M_{kj}}{\partial \theta_i} \dot{\theta}_k \dot{\theta}_j - \frac{\partial V}{\partial \theta_i}$$

Example : Two-Link Planar Manipulator



$$g_{c1}(\theta) = \begin{Bmatrix} \frac{1}{2}l_1 \cos(\theta_1) \\ \frac{1}{2}l_1 \sin(\theta_1) \\ \theta_1 \end{Bmatrix}$$

$$g_{c2}(\theta) = \begin{Bmatrix} l_1 \cos(\theta_1) + \frac{3}{5}l_2 \cos(\theta_1 + \theta_2) \\ l_1 \sin(\theta_1) + \frac{3}{5}l_2 \sin(\theta_1 + \theta_2) \\ \theta_1 + \theta_2 \end{Bmatrix}$$

\Rightarrow

$$\overset{\text{body}}{J}_{c1}(\theta) = \begin{bmatrix} 0 & 0 \\ \frac{1}{2}l_1 & 0 \\ 0 & 0 \end{bmatrix}$$

$$\overset{\text{body}}{J}_{c2}(\theta) = \begin{bmatrix} 0 & l_1 \sin(\theta_2) \\ 0 & l_1 \cos(\theta_2) + \frac{3}{5}l_2 \\ 1 & 1 \end{bmatrix}$$

$$L(\theta, \dot{\theta}) = T(\theta, \dot{\theta})$$

• no potential energy

$$T = \frac{1}{2} (\xi_1^{\text{body}})^T M_1 \xi_1^{\text{body}} + \frac{1}{2} (\xi_2^{\text{body}})^T M_2 \xi_2^{\text{body}}$$

$$\text{where } M_1 = \begin{bmatrix} m_1 & 0 & 0 \\ 0 & m_1 & 0 \\ 0 & 0 & I_1 \end{bmatrix} \quad \nparallel \quad M_2 = \begin{bmatrix} m_2 & 0 & 0 \\ 0 & m_2 & 0 \\ 0 & 0 & I_2 \end{bmatrix}$$

\Rightarrow

$$T(\theta, \dot{\theta}) = \frac{1}{2} \dot{\theta}^T (\overset{\text{body}}{J}_{c1}(\theta))^T M_1 \overset{\text{body}}{J}_{c1}(\theta) \dot{\theta} + \frac{1}{2} \dot{\theta}^T (\overset{\text{body}}{J}_{c2}(\theta))^T M_2 \overset{\text{body}}{J}_{c2}(\theta) \dot{\theta}$$

$$\begin{aligned}
& \frac{1}{2} \dot{\theta}^T (\mathbf{J}_{c_1}^{body}(\theta))^T M_1 \mathbf{J}_{c_1}^{body}(\theta) \dot{\theta} \\
= & \frac{1}{2} \dot{\theta}^T \begin{bmatrix} 0 & \frac{1}{2}l_1 & 1 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} m_1 & 0 & 0 \\ 0 & m_1 & 0 \\ 0 & 0 & I_1 \end{bmatrix} \begin{bmatrix} 0 & 0 \\ \frac{1}{2}l_1 & 0 \\ 1 & 0 \end{bmatrix} \dot{\theta} \\
= & \frac{1}{2} \dot{\theta}^T \begin{bmatrix} 0 & \frac{1}{2}l_1 & 1 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} 0 & 0 \\ \frac{1}{2}m_1 l_1 & 0 \\ I_1 & 0 \end{bmatrix} \dot{\theta} \\
= & \frac{1}{2} \dot{\theta}^T \begin{bmatrix} \frac{1}{4}m_1 l_1^2 + I_1 & 0 \\ 0 & 0 \end{bmatrix} \dot{\theta}
\end{aligned}$$

$$\begin{aligned}
& \frac{1}{2} \dot{\theta}^T (\mathbf{J}_{c_2}^{body}(\theta))^T M_2 \mathbf{J}_{c_2}^{body}(\theta) \dot{\theta} \\
= & \frac{1}{2} \dot{\theta}^T \begin{bmatrix} l_1 \sin(\theta_2) & l_1 \cos(\theta_2) + \frac{3}{5}l_2 & 1 \\ 0 & \frac{3}{5}l_2 & 1 \end{bmatrix} \begin{bmatrix} m_2 & 0 & 0 \\ 0 & m_2 & 0 \\ 0 & 0 & I_2 \end{bmatrix} \begin{bmatrix} l_1 \sin(\theta_2) & 0 \\ l_1 \cos(\theta_2) + \frac{3}{5}l_2 & \frac{3}{5}l_2 \\ 1 & 1 \end{bmatrix} \dot{\theta} \\
= & \frac{1}{2} \dot{\theta}^T \begin{bmatrix} l_1 \sin(\theta_2) & l_1 \cos(\theta_2) + \frac{3}{5}l_2 & 1 \\ 0 & \frac{3}{5}l_2 & 1 \end{bmatrix} \begin{bmatrix} m_2 l_1 \sin(\theta_2) & 0 \\ m_2 (l_1 \cos(\theta_2) + \frac{3}{5}l_2) & \frac{3}{5}m_2 l_2 \\ I_2 & I_2 \end{bmatrix} \dot{\theta} \\
= & \frac{1}{2} \dot{\theta}^T \begin{bmatrix} m_2 l_1^2 \sin^2(\theta_2) + m_2 (l_1 \cos(\theta_2) + \frac{3}{5}l_2)^2 + I_2 & \frac{3}{5}m_2 l_2 (l_1 \cos(\theta_2) + \frac{3}{5}l_2) + I_2 \\ \frac{3}{5}m_2 l_2 (l_1 \cos(\theta_2) + \frac{3}{5}l_2) + I_2 & \frac{9}{25}m_2 l_2^2 + I_2 \end{bmatrix} \dot{\theta}
\end{aligned}$$

$$\Rightarrow = \frac{1}{2} \dot{\theta}^T \begin{bmatrix} m_2 l_1^2 + \frac{9}{25}l_2^2 + 2 \cdot \frac{3}{5}m_2 l_1 l_2 \cos(\theta_2) + I_2 & \frac{3}{5}m_2 l_2 (l_1 \cos(\theta_2) + \frac{3}{5}l_2) + I_2 \\ \frac{3}{5}m_2 l_2 (l_1 \cos(\theta_2) + \frac{3}{5}l_2) + I_2 & \frac{9}{25}m_2 l_2^2 + I_2 \end{bmatrix} \dot{\theta}$$

~~(1)~~

$$L(\theta_1, \dot{\theta}_1) = \frac{1}{2} \dot{\theta}_1^2 \left[\begin{array}{cc} \frac{1}{4} m_1 l_1^2 + M_2 l_1^2 + \frac{9}{25} m_2 l_2^2 + \frac{6}{5} M_2 l_1 l_2 \cos(\theta_2) + I_1 + I_2 & \frac{3}{5} m_2 l_2 (l_1 \cos(\theta_2) + \frac{3}{5} l_2) + I_2 \\ \frac{3}{5} m_2 l_2 (l_1 \cos(\theta_2) + \frac{3}{5} l_2) + I_2 & \frac{9}{25} m_2 l_2^2 + I_2 \end{array} \right] \dot{\theta}_1$$

This is M(0)

now, what about the equations of motion?

just need to figure out:

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{\theta}} - \frac{\partial L}{\partial \theta} = \tau$$

\sum torques applied to joint actuators.

it will be of the form

$$M(\theta) \ddot{\theta} + C(\theta, \dot{\theta}) \dot{\theta} + N(\theta, \dot{\theta}) = \Theta \tau$$

[we don't have this part.]

$$C(\theta, \dot{\theta}) = \frac{1}{2} \sum_{k=1}^n \left(\frac{\partial M_{ij}}{\partial \theta_k} + \frac{\partial M_{ik}}{\partial \theta_j} - \frac{\partial M_{kj}}{\partial \theta_i} \right) \dot{\theta}_k$$

$$\frac{\partial M}{\partial \theta_1} = \begin{bmatrix} 0 & 0 \\ 0 & 0 \end{bmatrix} \quad \frac{\partial M}{\partial \theta_2} = \begin{bmatrix} -\frac{6}{5}m_2l_1l_2\sin(\theta_2) & -\frac{3}{5}m_2l_1l_2\sin(\theta_2) \\ -\frac{3}{5}m_2l_1l_2\sin(\theta_2) & 0 \end{bmatrix}$$

$$C_{ijL}(\theta, \dot{\theta}) = \begin{bmatrix} 0 & -\frac{6}{5}M_2l_1l_2\sin(\theta_2)\dot{\theta}_1 \\ \frac{6}{5}M_2l_1l_2\sin(\theta_2)\dot{\theta}_1 & 0 \end{bmatrix}$$

$$C_{ij2}(\theta, \dot{\theta}) = \begin{bmatrix} -\frac{6}{5}m_2l_1l_2\sin(\theta_2)\dot{\theta}_2 & -\frac{6}{5}m_2l_1l_2\sin(\theta_2)\dot{\theta}_2 \\ 0 & 0 \end{bmatrix}$$

\Rightarrow

$$C(\theta, \dot{\theta}) = \begin{bmatrix} -\frac{6}{5}m_2l_1l_2\sin(\theta_2)\dot{\theta}_2 & -\frac{6}{5}m_2l_1l_2\sin(\theta_2)(\dot{\theta}_1 + \dot{\theta}_2) \\ \frac{6}{5}m_2l_1l_2\sin(\theta_2)\dot{\theta}_1 & 0 \end{bmatrix}$$

which can be used to obtain the equations of motion as per

$$M(\theta)\ddot{\theta} + C(\theta, \dot{\theta})\dot{\theta} = \tau.$$