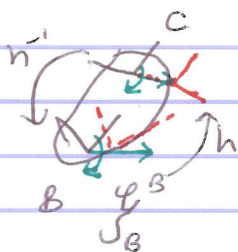


~~It~~ With respect to frame B to configuration B'. It does so under some constant body velocity

$$\mathcal{G}_B^B = \begin{Bmatrix} v \\ \omega \end{Bmatrix} \quad \begin{array}{l} v = \text{constant} \\ \omega = \text{constant} \end{array}$$

ok, but what does the reference frame C experience?

$$\mathcal{G}_C^C = h^{-1} \mathcal{G}_B^B h$$



$$h = h_C^B \Rightarrow \mathcal{G}_C^C = h_C^B \mathcal{G}_B^B h_C^B$$

In homogeneous coordinates:

$$\hat{\mathcal{G}}_C^C = h^{-1} \hat{\mathcal{G}}_B^B h = \begin{bmatrix} R_h & dh \\ 0 & 1 \end{bmatrix}^{-1} \begin{bmatrix} \hat{\omega} & v \\ 0 & 0 \end{bmatrix} \begin{bmatrix} R_h & dh \\ 0 & 1 \end{bmatrix}$$

$$= \begin{bmatrix} R_h^T & -R_h^T dh \\ 0 & 1 \end{bmatrix} \begin{bmatrix} \hat{\omega} & v \\ 0 & 0 \end{bmatrix} \begin{bmatrix} R_h & dh \\ 0 & 1 \end{bmatrix}$$

$$= \begin{bmatrix} R_h^T \hat{\omega} & R_h^T v \\ 0 & 0 \end{bmatrix} \begin{bmatrix} R_h & dh \\ 0 & 1 \end{bmatrix}$$

$$= \begin{bmatrix} R_h^T \hat{\omega} R_h & R_h^T \hat{\omega} dh + R_h^T v \\ 0 & 0 \end{bmatrix}$$

$$\begin{aligned} R_h^T \hat{\omega} R_h &= -\omega R_h^T J R_h \\ &= -\omega J = \hat{\omega} \end{aligned}$$

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$$\mathcal{G}_c^c = \begin{bmatrix} \hat{\omega} & 1 & R_h^T \hat{\omega} d_h + R_h^T v \\ 0 & -1 & 0 \end{bmatrix}$$

⇒ so, homogeneous coordinates, we can just ~~multiply~~ do matrix multiplication, what if I wanted a formula for the adjoint in vector form?

$$\mathcal{G}_c^c = \begin{Bmatrix} v_c^c \\ \omega_c^c \end{Bmatrix} = \text{Ad}_{h^{-1}} \begin{Bmatrix} v_b^b \\ \omega_b^b \end{Bmatrix}$$

First I want to know what $\text{Ad}_h \mathcal{G}$ is, then worry about the inverse.

$$\text{Ad}_h \hat{\mathcal{G}} = h \hat{\mathcal{G}} h^{-1} \quad \text{where } \mathcal{G} = \begin{Bmatrix} v \\ \omega \end{Bmatrix}$$

$$= \begin{bmatrix} R_h & d_h \\ 0 & 1 \end{bmatrix} \begin{bmatrix} \hat{\omega} & v \\ 0 & 0 \end{bmatrix} \begin{bmatrix} R_h^T & -R_h^T d_h \\ 0 & 1 \end{bmatrix}$$

$$= \begin{bmatrix} R_h \hat{\omega} R_h^T & -R \hat{\omega} R^T d_h + R v \\ 0 & 0 \end{bmatrix}$$

$$= \begin{bmatrix} \hat{\omega} & -\hat{\omega} d_h + R v \\ 0 & 0 \end{bmatrix}$$

My interest is vector form so I want

$$(\text{Ad}_h \hat{\mathcal{G}})^v = \begin{Bmatrix} -\hat{\omega} d_h + R v \\ \omega \end{Bmatrix}$$

Can I get $(\text{Ad}_h \hat{\mathcal{G}}) = \text{MATRIX} \cdot \mathcal{G}$



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$$= \begin{Bmatrix} \omega^T J^T d_h + R_h^T v \\ \omega \end{Bmatrix} = \begin{bmatrix} R & J^T d_h \\ 0 & 1 \end{bmatrix} \begin{Bmatrix} v \\ \omega \end{Bmatrix}$$

So, in vector form for y ,

$$\star \text{Ad}_h y = \begin{bmatrix} R_h & J^T d_h \\ 0 & 1 \end{bmatrix} \begin{Bmatrix} v \\ \omega \end{Bmatrix} \quad \text{for } y = \begin{Bmatrix} v \\ \omega \end{Bmatrix}$$

↳ Back to our story

$$\text{Ad}_{h^{-1}} y_B = \begin{bmatrix} R_h^T & -J^T R_h^T d_h \\ 0 & 1 \end{bmatrix} \begin{Bmatrix} v \\ \omega \end{Bmatrix} = \begin{Bmatrix} R_h^T v - J^T R_h^T \omega d_h \\ \omega \end{Bmatrix}$$

\downarrow
 $h^{-1} = \begin{bmatrix} R_h^T & -R_h^T d_h \\ 0 & 1 \end{bmatrix}$

Almost agrees, except for ordering of J^T and R_h^T but,

$$R_h J^T R_h^T = J^T$$

\Rightarrow mult. by R_h^T on the left

$$\cancel{R_h^T R_h} J^T R_h^T = \cancel{R_h^T} J^T$$

$$J^T R_h^T = R_h^T J^T$$

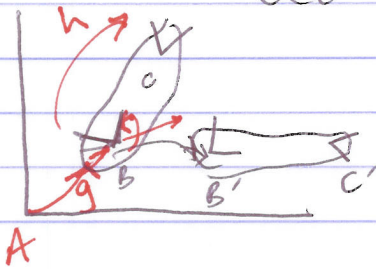
so, it actually agrees.

SYNOPSIS:

$$\text{Ad}_h y = \begin{bmatrix} R_h & J^T d_h \\ 0 & 1 \end{bmatrix} y \quad \text{in vector form } y = \begin{Bmatrix} v \\ \omega \end{Bmatrix}$$

OR

$$\text{Ad}_h \hat{y} = h \hat{y} h^{-1} \quad \text{in homogeneous form.}$$



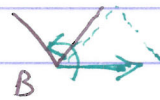
Body velocity: \mathcal{E}_B^B
 ↑
 velocity of body w.r.t
 coordinate frame of the
 body.

Suppose, I would rather know the velocity of another point on the body whose relative configuration is h (from B to C)

$$\mathcal{E}_C^C = \text{Ad}_{h^{-1}} \mathcal{E}_B^B \quad \text{where } h = h_C^B$$

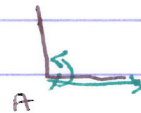
Adjoint didn't change rotational velocity.
 For SEC2) (e.g. planar case), rotations do not affect angular velocity.

$$\mathcal{E}_B^A = g_B^A * \mathcal{E}_B^B$$



$$\mathcal{E}_B^B = \begin{bmatrix} 1 \\ -1 \\ 1 \end{bmatrix}$$

$$= \begin{bmatrix} B & 1 & 0 \\ 0 & 1 & 1 \end{bmatrix} \begin{bmatrix} v \\ w \end{bmatrix}$$



$$\mathcal{E}_B^A = \begin{bmatrix} \sqrt{2} \\ 0 \\ 1 \end{bmatrix}$$

$$\mathcal{E}_B^A = \begin{bmatrix} Rv \\ w \end{bmatrix}^A$$

How to change
reference frames of a
vector's coordinate
representation.

$$R\vec{v} + \vec{w}$$

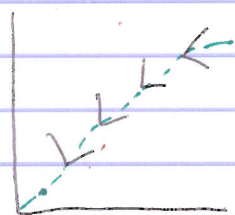
$$= R \begin{bmatrix} v_x \\ v_y \end{bmatrix} + \begin{bmatrix} w_x \\ w_y \end{bmatrix}$$

$$\text{Ad}_h = \begin{bmatrix} R_h & J^T \\ 0 & 1 \end{bmatrix} \begin{bmatrix} v \\ w \end{bmatrix}$$

Why do I need this?

Because we may be given \mathcal{E}_B^A and we really want \mathcal{E}_B^B

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$$g(t) = g_B^A(t)$$

$$\Rightarrow \frac{d}{dt}$$

$$\dot{g}_B^A(t)$$

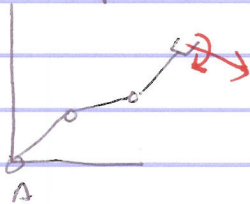
← this is $\xi_B^A(t)$

How did I get into body frame?

We did $\xi_B^B(t) = g_A^B(t) * \xi_B^A(t)$ (on page 8-1)

$$= (g_B^A(t))^{-1} * \xi_B^A(t)$$

for manipulators



$$g_B^A = g_B^A(\theta_1, \dots, \theta_n)$$

Spatial Velocity:

defined as:

$$\begin{aligned} \dot{g}(t) g^{-1}(t) &= \left\{ \dot{g}_B^A(t) (g_B^A(t))^{-1} \right. \\ &= \dot{g}_B^A(t) g_A^B(t) \\ &= \xi_A^A(t) \end{aligned}$$

what is this?

what?!

What if I do the following:

$$\text{Ad}_{(g_B^A)^{-1}} \xi_B^B = (g_B^A) \xi_B^B (g_B^A)^{-1} =$$

$$\underbrace{g_A^B \xi_B^B}_{\xi_A^A} = \cancel{g_B^A g_A^B} \dot{g}_B^A g_A^B$$

The spatial velocity is the velocity associated with the observer frame (A) moving as though it were rigidly attached to the body frame (B), which itself was moving.

$$\xi_A^A = \text{Ad}_{g^{-1}} \xi_B^B \quad \text{where } g^* = g_B^A$$

So, what is it for SE(2)?

$$\dot{g}_B^A(t) (g_B^A(t))^{-1} = \begin{bmatrix} \frac{dR(\theta(t))}{dt} \dot{\theta}(t) & \dot{d}(t) \\ 0 & 0 \end{bmatrix} \begin{bmatrix} R^T(\theta(t)) & 0 \\ 0 & 1 \end{bmatrix}$$

$$\text{for } g_B^A(t) = \begin{bmatrix} R(\theta(t)) & d(t) \\ 0 & 1 \end{bmatrix}$$

$$= \begin{bmatrix} \frac{dR(\theta(t))}{dt} R^T(\theta(t)) \dot{\theta}(t) & \dot{d}(t) \\ 0 & 0 \end{bmatrix}$$

$$= \begin{bmatrix} -J^T \dot{\theta} & J^T \dot{d} + \dot{d} \\ 0 & 0 \end{bmatrix}$$

$$\hat{\xi}_A^A(t) = \begin{bmatrix} \hat{\theta} & -\hat{\theta} d + \dot{d} \\ 0 & 0 \end{bmatrix}$$

Homogeneous form of spatial velocity
velocity of frame A moving in a manner consistent
with being rigidly attached to the moving frame B.

09/14/07

ECE 4560

Vectors of $SE(2)$. Discussed:

1) Linear structure

2) Adjoint & change of reference frame for a moving body

TO DO: 3) Relationship b/w elements and vectors (velocities)
Mapping of displacement to vector and vice-versa

3) Exponential Representation

Notation:

a configuration is given by a Lie group element,

$$g \in SE(2)$$

↳ special Euclidean

a vector associated to $SE(2)$ is given by aLie algebra element, $\xi \in \mathfrak{se}(2)$

We want the equivalent of

$$\dot{x} = v$$

for g and ξ In $SE(2)$, ξ looks like $\xi = \begin{Bmatrix} v \\ \omega \end{Bmatrix}$

$$\text{Given } \dot{g} = g * \xi$$

what happens if $\omega = 0$

$$\Rightarrow \dot{g} = \begin{bmatrix} R & 0 \\ 0 & 1 \end{bmatrix} \begin{Bmatrix} v \\ \omega \end{Bmatrix} = \begin{Bmatrix} Rv \\ 0 \end{Bmatrix}$$

$$\Rightarrow \dot{x} = R_{11} v_x + R_{12} v_y = \omega_x$$

$$\dot{y} = R_{21} v_x + R_{22} v_y = \omega_y$$

$$\dot{\theta} = 0$$

$$R = \begin{bmatrix} R_{11} & R_{12} \\ R_{21} & R_{22} \end{bmatrix}$$

$$\Rightarrow \begin{aligned} \dot{x} &= \omega x \\ \dot{y} &= \omega y \\ \dot{\theta} &= 0 \end{aligned}$$

What happens when $\omega \neq 0$?
In this case,

$$\begin{aligned} \dot{x} &= R_{11}(\theta)v_x + R_{12}(\theta)v_y \\ \dot{y} &= R_{21}(\theta)v_x + R_{22}(\theta)v_y \\ \dot{\theta} &= \omega \end{aligned}$$

$$\Rightarrow \begin{aligned} \dot{x} &= v_x \cos \theta - v_y \sin \theta \\ \dot{y} &= +v_x \sin \theta + v_y \cos \theta \\ \dot{\theta} &= \omega \end{aligned} \quad \begin{array}{l} \text{Is there a solution} \\ \text{to this?} \end{array}$$

Justify my solution:
Recall: $\dot{g} = g \cdot \xi$

In linear control, you guys worked with:

$$\dot{x} = Ax$$

what was the integral?

Solution was an exponential,

$$x(t) = e^{At} x_0$$

Basic definition:
$$e^{At} = \sum_{n=1}^{\infty} \frac{(At)^n}{n!}$$

Best approach is to diagonalize:

$$A = T \Sigma T^{-1} = A_d T \Sigma$$

$$e^A = e^{A_d T \Sigma} = A_d e^{\Sigma} T^{-1}$$

$$e^{\begin{bmatrix} \sigma_1 & & 0 \\ & \ddots & \\ 0 & & \sigma_n \end{bmatrix}} = \begin{bmatrix} e^{\sigma_1} & & 0 \\ & \ddots & \\ 0 & & e^{\sigma_n} \end{bmatrix}$$

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In the end I get a transition matrix Φ , $X_{\text{new}} = \Phi X_{\text{old}}$

for elements of $SE(2)$,

there is an exponential also.

It is defined:

$$\exp(\hat{\xi}\tau) = \begin{cases} \text{if } \omega = 0 & \text{then it is } \begin{bmatrix} 1 & | & v\tau \\ 0 & | & 1 \end{bmatrix} \\ \text{if } \omega \neq 0 & \text{then it is } \begin{bmatrix} e^{\hat{\omega}\tau} & | & -\frac{1}{\omega}(1 - e^{\hat{\omega}\tau})Jv \\ 0 & | & 1 \end{bmatrix} \end{cases}$$

$$\exp: SE(2) \times \mathbb{R} \rightarrow SE(2)$$

The inverse operation is called the logarithm, and will be ~~denoted~~ denoted by \ln .

It has one of two forms:

$$\ln: SE(2) \rightarrow SE(2) \quad (\text{time not specified})$$

OR

$$\ln: SE(2) \rightarrow SE(2) \times \mathbb{R} \quad (\text{time specified})$$

Given $g \in SE(2)$, then

$$g = \begin{bmatrix} R & | & d \\ 0 & | & 1 \end{bmatrix}$$

$$(\hat{\xi}, \tau) = \ln g = \begin{cases} \text{if } R = \mathbb{1}, \text{ then } v = \frac{d}{\|d\|} \text{ \& \# \& \# } \tau = \|d\| \\ \text{else,} \end{cases}$$

$$\tau = 2 \tan^{-1} \left(\frac{\|R_{21}\|}{R_{11}} \right)$$

$$\omega = 1$$

$$v = J(1 - R)^{-1} d$$

The alternative is

$$\xi = \ln g = \begin{cases} \bar{y} & R = \mathbb{1}, \text{ then} \\ \text{else} \end{cases}$$

$$\begin{aligned} v &= \bar{d} \\ w &= 0 \end{aligned}$$

$$\begin{aligned} w &= 2 \tan^{-1}(R_{21}, R_{11}) \\ v &= w^T (1 - R)^{-1} d \end{aligned}$$